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DOI: <https://doi.org/10.1093/mnras/stv2582>

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ZORA URL: <https://doi.org/10.5167/uzh-129992>

Journal Article

Originally published at:

Bruderer, Claudio; Read, Justin I; Coles, Jonathan P; Leier, Dominik; Falco, Emilio E; Ferreras, Ignacio; Saha, Prasenjit (2016). Light versus dark in strong-lens galaxies: dark matter haloes that are rounder than their stars. *Monthly Notices of the Royal Astronomical Society*, 456(1):870-884.

DOI: <https://doi.org/10.1093/mnras/stv2582>

Light versus dark in strong-lens galaxies: Dark matter haloes that are rounder than their stars

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5 November 2015

ABSTRACT

We measure the projected density profile, shape and alignment of the stellar and dark matter mass distribution in 11 strong-lens galaxies. We find that the projected dark matter density profile – under the assumption of a Chabrier stellar initial mass function – shows significant variation from galaxy to galaxy. Those with an outermost image beyond ~ 10 kpc are very well fit by a projected NFW profile; those with images within 10 kpc appear to be more concentrated than NFW, as expected if their dark haloes contract due to baryonic cooling. We find that over several half-light radii, the dark matter haloes of these lenses are rounder than their stellar mass distributions. While the haloes are never more elliptical than $e_{dm} = 0.2$, their stars can extend to $e_* > 0.2$. Galaxies with high dark matter ellipticity and weak external shear show strong alignment between light and dark; those with strong shear ($\gamma \gtrsim 0.1$) can be highly misaligned. This is reassuring since isolated misaligned galaxies are expected to be unstable. Our results provide a new constraint on galaxy formation models. For a given cosmology, these must explain the origin of both very round dark matter haloes and misaligned strong-lens systems.

Key words: Gravitational lensing: strong – galaxies: structure – galaxies: haloes – galaxies: formation – galaxies: elliptical and lenticular, cD.

1 INTRODUCTION

The ellipticity and shape of the stellar component relative to their host dark matter halo encodes information both about our current cosmological model Λ CDM and galaxy formation (e.g. Dubinski 1994; Ibata et al. 2001; Kazantzidis et al. 2004; Macciò et al. 2007; Debattista et al. 2008; Lux et al. 2012; Read 2014). ‘Dark-matter-only’ (DMO) simulations in Λ CDM predict dark matter haloes that are triaxial (Dubinski & Carlberg 1991; Warren et al. 1992; Navarro et al. 1996; Jing & Suto 2002), with mean ‘shape parameter’ $\langle q \rangle = (b + c)/2a \sim 0.8$ (where $a > b > c$ are the long, intermediate and short axes of the figure; Macciò et al. 2007). This corresponds to a typically *prolate* halo. However, including ‘baryons’ (stars and gas) in the models pro-

duces haloes that are significantly rounder and – at least for disc galaxies – well-aligned with the light distribution (Katz & Gunn 1991; Dubinski 1994; Debattista et al. 2008). Halo shapes and alignments also constrain alternative gravity models (Mortlock & Turner 2001; Helmi 2004; Read & Moore 2005; Ferreras et al. 2012; Debattista et al. 2013). If stars dominate the mass of the galaxy, we expect the light and mass distribution to be highly correlated; if dark matter is present, however, such correlations can, at least in principle, be broken.

Strong lensing provides a unique probe of the alignment and shape of the total mass distribution in galaxies (e.g. Blandford & Narayan 1986; Schneider et al. 1992; Keeton et al. 1998; Kochanek et al. 2000; Koopmans et al. 2006; Auger et al. 2007; Ferreras et al. 2008; Auger et al. 2010; Leier et al. 2012). For ‘red and dead’ ellipticals that are largely devoid of gas, their baryonic content can be mapped

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through stellar population synthesis modelling of their light distribution alone (e.g. Ferreras et al. 2005; Treu et al. 2006; Ferreras et al. 2008). Furthermore, such systems are dense enough to produce strong lensing effects, opening up the possibility of directly comparing the light and mass in these galaxies (Keeton et al. 1998; Ferreras et al. 2008; Treu et al. 2009; Sluse et al. 2012). Previous work in the literature has found that the light and mass are well-aligned (though a mis-match of up to 10° is not uncommon; e.g. Sluse et al. 2012). However, results on the ellipticity of light and mass agree less well, with Sluse et al. (2012) finding a strong correlation and Keeton et al. (1998) and Ferreras et al. (2008) finding none. It is difficult, however, to compare the results between these different studies because they use different lens modelling techniques; different definitions of ellipticity; and different radii over which the shapes and alignments are probed. Furthermore, none to date have applied their methodology to mock data to determine the robustness of the results.

Weak lensing can also be used to probe the shape and alignment between the luminous and dark matter distributions within galaxies. However, this requires that the galaxies are ‘stacked’ (Brainerd & Wright 2000; Natarajan & Refregier 2000), typically by aligning the major axes of the light distribution (e.g. Hoekstra et al. 2004). This will lead, in absence of contaminating effects (e.g. intrinsic alignments; Hirata & Seljak 2004), to an isotropic mean shear around galaxies if the mass distributions are spherical, or if dark and light are randomly aligned. The weak lensing measurement has been performed on data by various groups (Hoekstra et al. 2004; Mandelbaum et al. 2006; Parker et al. 2007; van Uitert et al. 2012) with inconclusive results. The analyses differ mainly in the data sets used; the selection of lens galaxies; the estimators employed; and the treatment of systematics (see e.g. Schrabback et al. 2015). For example, Clampitt & Jain (2015) have recently made the measurement on Luminous Red Galaxies (LRG) taken from the Sloan Digital Sky Survey (SDSS) reporting a significant detection of dark matter halo ellipticity; while Schrabback et al. (2015) performed a similar measurement on blue and red galaxies from the Canada France Hawaii Lensing Survey (CFHTLenS) finding no definite detection.

Recently, Coles et al. (2014) introduced a new non-parametric lens tool, GLASS. Applying this to a large suite of mock data, we showed that mass and light can only reliably be disentangled in strong lens systems if: i) there are at least four images; and ii) time delay data are available and/or the stellar mass contributes significantly to the potential. In this paper, we collate data of the above quality, compiling a sample of 11 strong lens galaxies. We apply GLASS to these lenses to non-parametrically measure the shape and alignment of the stars and *dark matter* in these lens galaxies, for the first time. This differs from previous works that have all compared the light distribution with the total mass, rather than the dark matter. Since the stellar component often dominates the central potential, the total mass naturally correlates with the light, potentially masking theoretically interesting results about the dark matter distribution. Our comparison between the stellar and dark matter components is made possible by the fact that GLASS uses the stellar mass distribution as a prior on the mass map, ensuring that the dark matter map is always positive.

This paper is organised as follows. In Section 2, we briefly review the GLASS code and define our method to assess shape and alignment of lens galaxies. In Section 3, we present our data compilation with references. In Section 4, we present our results. We discuss the implications of these results in Section 5. Finally, we conclude in Section 6.

Throughout this work we assume a flat Λ CDM model with matter density $\Omega_m = 0.28$, dark energy density $\Omega_\Lambda = 0.72$, and inverse Hubble constant $H_0^{-1} = 13.7$ Gyr.

2 THE LENS MODELS

While it is possible to compute shape parameters for a lensing galaxy by fitting a parametric shape to the lens, it makes the definition of a shape estimate dependent on the parametric form being used. Moreover, the commonly-used parametric forms for modelling lensing galaxies (e.g. Keeton 2001) do not allow for features like twisting isodensity contours, which can arise from the projection of triaxial ellipsoids with no intrinsic twists (e.g. Binney 1978, and references therein). To avoid these problems, we use non-parametric ellipticity estimators, defined as follows.

Lenses are modelled as free-form mass distributions, consisting of mass tiles or pixels. Such lens models are also called non-parametric, which really just means that many more parameters than data constraints are used. Thus the lens models are non-unique, and hence we build *ensembles* of these free-form mass distributions. From such a mass map, an inertia tensor is defined as:

$$I = \begin{pmatrix} \sum_{\theta} M(\theta) \theta_y^2 & -\sum_{\theta} M(\theta) \theta_x \theta_y \\ -\sum_{\theta} M(\theta) \theta_x \theta_y & \sum_{\theta} M(\theta) \theta_x^2 \end{pmatrix} \quad (1)$$

where the sum is over mass pixels, and $M(\theta)$ is the mass in a pixel. The eigenvectors of this inertia tensor give the ellipticity axes, and our ellipticity estimator is

$$e \equiv \frac{\lambda_1 - \lambda_2}{\lambda_1 + \lambda_2}, \quad (2)$$

where λ_1 and λ_2 ($\lambda_1 \geq \lambda_2$) are the eigenvalues. We use this ellipticity estimator to quantify the shapes both of the stellar-mass distribution e_* and of the dark-matter distribution e_{dm} . By default, we use the dark and stellar mass distributions out to the outermost lensing image. In Figure 3, we explore the effect of averaging instead over multiples of the half-light radius R_e . The mass pixels used in this work are larger than image pixels, but at most 4% of the diameter of the full mass map. The central region has smaller pixels, to allow for a density cusp at the centre.

We estimate the orientation of the distributions of luminous and dark matter by computing the angles θ_* and θ_{dm} of the eigenvector corresponding to the eigenvalue λ_1 ($\lambda_1 \geq \lambda_2$) relative to the x -axis. The misalignment angle $\Delta\theta$ between the distributions is then defined as:

$$\Delta\theta = \theta_{dm} - \theta_*. \quad (3)$$

Each model ensemble consists of 10^4 models. We apply our shape and alignment estimators to each model, in this way obtaining the median and 68% confidence intervals of e and $\Delta\theta$, for each lens.

The non-parametric mass models themselves are constructed using the GLASS framework for modelling multiple-image lenses (Coles et al. 2014). GLASS produces ensembles

Lens	z_L	z_S	$\Delta\phi$ [kpc]	R_L/R_e	$M_s(< 2R_e)$ [$10^{10} M_\odot$]	Env	θ_{Env} [°]
0047	0.485	3.60	12.82	1.45 ± 0.04	11.58 ± 0.43	G(9) ^a	...
0414	0.960	2.64	16.01	1.85 ± 0.05	19.90 ± 2.29
0712	0.410	1.34	6.82	1.15 ± 0.03	5.46 ± 0.52
0911	0.769	2.8	23.16	3.09 ± 0.05	14.52 ± 1.87	C ^b	-160 ^b
0957	0.356	1.41	29.98	3.51 ± 0.04	20.92 ± 0.96	C ^c	60 ^f
1115	0.310	1.72	10.76	2.86 ± 0.06	5.61 ± 1.19	G(13) ^a	-125 ^g
1422	0.337	3.62	6.02	4.49 ± 0.06	2.91 ± 0.45	G(17) ^a	145 ^g
1608	0.630	1.39	13.92	1.82 ± 0.01	27.99 ± 1.63	G(8) ^d	30 resp. -5 ^{d,h}
2016	1.010	3.3	26.22	6.12 ± 0.14	6.34 ± 1.99	C(69) ^e	130 - 140 ^e
2045	0.870	1.28	14.46	1.48 ± 0.03	14.05 ± 1.03
2237	0.039	1.7	1.40	0.89 ± 0.01	1.15 ± 0.12

Table 1. The most relevant lens properties for this work are listed (for an expanded version of this table see Leier et al. 2011). References can be found in Section 3. The columns are from left to right: the lens redshift z_L ; the redshift of the source z_S ; the maximum angular separation between two images $\Delta\theta$; the ratio of the radius of the outermost image R_L and the effective Petrosian half-light radius of the lens R_e ; the stellar mass M_s within $2R_e$; information on the environment of the lens; and the position angle to the centroid of the corresponding group/cluster (measured north through east). C and G denote a known cluster or group environment, respectively. The number in the brackets is the number of confirmed members. Ellipses indicate lens galaxies with no group members identified so far.

^a Wong et al. (2011); ^b Morgan et al. (2001); ^c e.g. Garrett et al. (1992); ^d Fassnacht et al. (2006); ^e Toft et al. (2003); ^f Chartas et al. (2002); ^g Grant et al. (2004).

^h The position angle is given for the unweighted resp. the luminosity-weighted centroid estimation.

of mass maps, each of which reproduces exactly the observed image positions, image parities and time delays (if measured). The lensing data must be in the form of point images, which could be quasars or simply point-like features in otherwise extended images, but this was not a problem for the lenses studied in this work. The estimated stellar mass (explained below in Section 3) is taken as a lower limit on the total mass. The additional mass that needs to be put on top, in order to reproduce the lensing observables, is interpreted as dark matter. The lensing observables and assumed non-negative dark matter by themselves leave the mass distribution under-determined. Hence, an additional prior is used, requiring the mass distribution to be centrally concentrated, with the same centre as the galaxy light (but not necessarily the same shape). The precise formulation of the prior is given in Section 3 of Coles et al. (2014) (alternatively, for a more intuitive account of GLASS see Section 3.2 of Küng et al. 2015).

GLASS is closely related to an earlier code PIXELENS (Saha & Williams 2004; Coles 2008), but has a completely different code base. An important improvement is a better sampling algorithm for high dimensional spaces (Lubini & Coles 2012) that eliminates the excessive weight PIXELENS tended to give to extreme models.

Most relevant to the present work are Section 5.2 and Figure 8 of Coles et al. (2014), showing the recovery of λ_2/λ_1 . The typical errors in this quantity are 10–20%.

3 DATA

Stellar mass maps are taken from Leier et al. (2011). In that work, the surface brightness distribution of the galaxy images in all available HST bands was first fitted using GALFIT v2.03B (Peng et al. 2002). This was then converted to a stellar mass map using the stellar-population synthesis (SPS) models of Bruzual & Charlot (2003). The population synthesis is marginalised over the star formation epoch and time

scales and over the stellar metallicity. The initial mass function (IMF) was taken to be the log-normal form given by Chabrier (2003); a more bottom-heavy IMF would change the mass normalisation (cf. Schechter et al. 2014) but not the shape of the inferred stellar mass distribution (unless the IMF presents significant intrinsic deviations locally, see e.g. Martín-Navarro et al. 2015). We note that the effect of recent claims towards a non-universal IMF on stellar mass-to-light ratio (M/L) is still under debate. While for massive elliptical galaxies, dynamical studies seem to point at contributions at the level of the Salpeter IMF or slightly heavier (Cappellari et al. 2013), spectral line strength constraints can still accommodate a wider range of the stellar M/L, depending on the functional form of the IMF (Ferreras et al. 2013). For at least one lens in our sample, *Q2237+030*, the lensing estimates on the mass distribution are inconsistent with a Salpeter IMF (Ferreras et al. 2010). We additionally tested varying the IMF for one lens (*B1422+231*), finding that a Salpeter IMF is too heavy and gives no solutions. We will perform a more systematic study of the IMF and varying stellar M/L in a forthcoming publication.

The original sample we consider, is presented in Leier et al. (2011) and consists of 21 lens systems. It is a sub-sample of the CfA-Arizona Space Telescope LENS Survey¹ (CASTLES) sample. As shown in Coles et al. (2014), the recovery of the shape information of the lens is only good if there are at least four images. Therefore, we only consider the quad lenses and the twin double *Q0957+561*. This yields a sample consisting of 11 lenses. We discuss each of these systems, next. We note that all the position angles are measured north through east and relative to the lens galaxy.

3.1 The individual lenses

Q0047-2808 (hereafter *0047*) is a luminous early-type galaxy (Warren et al. 1996). It is part of a group of 9 mem-

¹ <http://www.cfa.harvard.edu/castles/>

Lens	Time delays [days]	Symmetric?	Comments
0047
0414
0712
0911	$\Delta t_{BA} = 146^{+4}_{-4}$	yes	...
0957	$\Delta t_{BA} = 416.5^{+1}_{-1}$	yes	...
1115	$\Delta t_{BA} = 12.0^{+2}_{-2}$
	$\Delta t_{DC} = 4.4^{+3.2}_{-2.4}$
1422	...	yes	...
1608	$\Delta t_{BA} = 31.5^{+2}_{-1.5}$
	$\Delta t_{CA} = 36.0^{+1.5}_{-1.5}$
	$\Delta t_{DA} = 77.0^{+2}_{-1}$
2016	Smaller pixels
2045	...	yes	$\gamma \leq 0.1$
2237

Table 2. Lenses where additional data or priors were used (cf. Section 2). Δt_{BA} and so on denote time delays between images as labelled in Table A1; γ refers to the allowed external shear; Symmetric ‘yes’ means the mass map is symmetric under rotation by 180° ; ‘Smaller pixels’ indicates that the mass pixels are 3% (rather than 4%) of the diameter of the whole mass map.

bers, all of which are spectroscopically confirmed (Wong et al. 2011).

MG0414+0534 (hereafter *0414*) is a passively evolving early-type galaxy (Tonry & Kochanek 1999). A close luminous satellite galaxy was found north-west of the lens (Schechter & Moore 1993). It, however, is more likely a foreground object with a redshift of 0.38 (Curran et al. 2011).

B0712+472 (hereafter *0712*) is an early-type galaxy (Jackson et al. 1998; Fassnacht & Cohen 1998). There is a foreground group of about 10 member galaxies at $z \sim 0.3$ (Fassnacht & Lubin 2002).

RXJ0911+0551 (hereafter *0911*) is an early-type galaxy (Bade et al. 1997; Sluse et al. 2012), and has measured time delays (Hjorth et al. 2002). It lies on the outskirts of a cluster of which X-ray emission can be detected yielding a temperature of 2.3 keV (Morgan et al. 2001). The center of the cluster lies at a position angle of about -160° . There is furthermore a satellite galaxy to the north-west (Kneib et al. 2000).

Q0957+561 (hereafter *0957*) is a cD galaxy lying close to the centre of a cluster with a high spiral galaxy-fraction (e.g. Garrett et al. 1992; Angonin-Willaime et al. 1994; Chartas et al. 1998). Due to the large image separation, large physical scales are probed. Identified X-ray emission furthermore locates the center of the cluster at a position angle of about 60° , close to the lens galaxy (Chartas et al. 2002). The lens has a measured time delay (e.g. Shalyapin et al. 2012). They however also find a three-day lag between the g- and r-bands, a disagreement at the 2σ -level. They argue that this effect can be accounted for by the presence of a substructure and chromatic dispersion. We find that the results do not change significantly for either estimate and choose therefore the g-band measurement.

PG1115+080 (hereafter *1115*) is an early-type galaxy (Weymann et al. 1980; Yoo et al. 2005). It has measured time delays (see e.g. Schechter et al. 1997). In this work, we use recent estimates of the time delays by Tsvetkova et al. (2010) that differ from previously found values by e.g. Barkana (1997). The environment of the lens was analysed

thoroughly (Momcheva et al. 2006; Wong et al. 2011). It is part of a small group of 13 members. Also, X-ray emission was detected from the corresponding group at a position angle of about -125° and it yields a temperature of 0.8 keV (Grant et al. 2004).

B1422+231 (hereafter *1422*) is an early-type galaxy (Patnaik et al. 1992; Yee & Ellingson 1994). Although time delays have been reported (Patnaik & Narasimha 2001), it is possible that the measurements are aliases of the actual values (Raychaudhury et al. 2003) and hence are not included in our analysis. The lens is part of a group with 16 spectroscopically confirmed member galaxies (Momcheva et al. 2006). Using newer data, additional members could be identified (Wong et al. 2011). Grant et al. (2004) also detect X-ray emission from the corresponding group originating from a position angle of about 145° at a temperature of 1.0 keV.

B1608+656 (hereafter *1608*) was first reported by Myers et al. (1995). It consists of two merging galaxies. The main galaxy is an early-type galaxy which is disrupted by a smaller, probably late-type galaxy (Surpi & Blandford 2003). The system has measured time delays (Fassnacht et al. 2002). The environment has been analysed and a group of 8 (9 if the merging galaxies are counted individually) member galaxies was identified (Fassnacht et al. 2006). Furthermore, the center was identified to be at a position angle of about 30° (resp. -5°) if the member galaxies’ positions are averaged without weighting (resp. weighted by luminosity). However, no significant X-ray emission was detected from the surrounding group (Dai & Kochanek 2005). The data seem to indicate four other groups along the line-of-sight.

MG2016+112 (hereafter *2016*) is a giant elliptical galaxy (Lawrence et al. 1984; Schneider et al. 1986). It is the farthest lens we consider in this sample. It lies in a cluster that consists of 69 probable galaxies (Toft et al. 2003). The cluster shows a high density of galaxies close to the lens in a south-east direction, at about $\sim 130 - 140^\circ$. Two lensed images are very close in this fold lens, but could be resolved by More et al. (2009).

B2045+265 (hereafter *2045*) is probably an elliptical galaxy (McKean et al. 2007). It was initially classified as a late-type Sa galaxy (Fassnacht et al. 1999), however the velocity dispersion seems too high. As the source redshift is rather low, a large lens mass is required. McKean et al. (2007) therefore conclude that it is more likely an elliptical galaxy. North-west of the lens there is evidence for a potential group at a similar redshift as the lens (Fassnacht et al. 1999). There is furthermore evidence for a dwarf satellite galaxy within the Einstein radius of the lens system (McKean et al. 2007). As the measurements are however inconclusive, we choose to not include this satellite in our models.

Q2237+030 (hereafter *2237*) is a barred spiral (Yee 1988). With a redshift of just $z \sim 0.04$ it is the closest lens system of the sample. Due to its low redshift, the probed physical scales are small, making the bulge component of the galaxy dominant. This makes the treatment of the stellar populations easier, as derivations of the stellar M/L are more difficult in the presence of star forming regions typical of disc galaxies, where there are no simple stellar populations usually assumed in SPS models.

Further information on the sample can be found in Leier et al. (2011) and Sluse et al. (2012). The most important quantities of each system are given in Table 1. Due to the

different angular separations of the images and the different angular diameter distances, in each lens galaxy we necessarily probe slightly different scales.

4 RESULTS

We have applied GLASS to the data set described in Section 3 consisting of 11 lenses. The reconstructed stellar and dark mass maps for each individual lens are given in the Appendix.

4.1 Radial dark matter profiles

Before discussing the shapes and alignments, we first measure the dark matter profiles of the strong lens galaxies, comparing these with Navarro-Frenk-White profiles (NFW; Navarro et al. 1996) given by (Klypin et al. 2011):

$$\frac{\rho(r)}{\rho_{crit}} = \frac{\delta_c}{(cr/r_{vir})(1 + cr/r_{vir})^2}, \quad (4)$$

where the critical density of the universe $\rho_{crit}(z)$ and $\delta_c(z)$ are parameters set by the cosmology considered, r is the radius from the centre, and c is the concentration parameter. The concentration is well-correlated with M_{vir} with some redshift dependence (Klypin et al. 2011):

$$c(M_{vir}, z) = c_0(z) \left(\frac{M_{vir}}{10^{12} h^{-1} M_\odot} \right)^{-0.075} \times \left[1 + \left(\frac{M_{vir}}{M_0(z)} \right)^{0.26} \right], \quad (5)$$

where $c_0(z)$ and $M_0(z)$ are fitting parameters (see Table 3 of Klypin et al. 2011). For each lens galaxy we choose the parameters at the redshift closest to the lens redshift.

For a circular NFW profile truncated at the virial radius r_{vir} , the projected surface mass density $\Sigma(x)$, where $x = cr/r_{vir}$, can be expressed as (Baltz et al. 2009):

$$\Sigma(x) = \frac{M_1}{r_{vir}^2} \frac{c^4}{2\pi(c^2 + 1)^2} \left\{ \frac{c^2 + 1}{x^2 - 1} [1 - F(x)] + 2F(x) - \frac{\pi}{\sqrt{c^2 + x^2}} + \frac{c^2 - 1}{c\sqrt{c^2 + x^2}} L(x) \right\}, \quad (6)$$

where

$$M_{vir} = M_1 \frac{\tau^2}{(\tau^2 + 1)^2} [(\tau^2 - 1)\ln\tau + \tau\pi - (\tau^2 + 1)] \quad (7)$$

and

$$F(x) = \frac{\cos^{-1}(1/x)}{\sqrt{x^2 - 1}} \quad \text{and} \quad L(x) = \ln \left(\frac{x}{\sqrt{c^2 + x^2} + c} \right). \quad (8)$$

The convergence $\kappa(x) = \Sigma(x)/\Sigma_{crit}$ can then be computed, where Σ_{crit} is the critical surface density at the lens' redshift. Finally, $\langle\kappa(R)\rangle$ is obtained by azimuthally averaging.

As the virial radius r_{vir} is related to the virial mass M_{vir} and the concentration parameter c is well-correlated with M_{vir} (Eq. 5), there is only one degree of freedom, M_{vir} . We calibrate the mass in order for the derived convergence profiles $\langle\kappa(R)\rangle$ of the reconstructed dark matter distribution and the NFW-profile to match at one pixel past the last image radius.

Figure 1 displays the results, while the virial masses (derived following the above procedure) are given in Table 3. In all cases the dark matter distribution either follows the NFW profile or is steeper, as expected if dark haloes contract in response to baryon cooling and star formation (e.g. Katz & Gunn 1991; Dubinski 1994; Debattista et al. 2008). At large radii, there is in all cases a sharp fall off in the GLASS dark matter density profiles. This owes to the edge of the mass map and for this reason we should not trust the reconstructed profiles outside of the outermost image. Interestingly, there is significant scatter from galaxy to galaxy in the shape of the azimuthally averaged dark matter density profile. This appears to correlate well with the physical radii probed. The galaxies that most closely follow an NFW profile – 0911, 0957, 0414, 1608, 2016, 2045 – all have an outermost image at radii > 10 kpc. The other lenses – 0047, 0712, 1115, 1422, 2237 – probe smaller radii and show signs of contraction with respect to NFW. In the most extreme example – 2237 – the outermost image probes only the bulge region of this galaxy (~ 0.8 kpc). The dark matter profile is substantially steeper than NFW while the derived virial mass is implausibly large (since we have not accounted for any such contraction in deriving it; see Table 3). Our results suggest that the dark matter haloes in these lensing galaxies are broadly consistent with contracted NFW profiles. Lenses with close-in images – that are expected to inhabit less massive dark matter haloes – show more apparent contraction. There could be several explanations for this behaviour. It could be a selection effect where lower mass galaxies can only strong lens if they are substantially contracted; it could owe to more massive haloes being less affected by baryons; or it could be that we simply cannot yet resolve the contraction at the centre of the more massive systems. We will explore such questions in a forthcoming paper where we will present the dark matter radial profiles of a larger sample of lenses, considering varying stellar initial mass function, and explicitly modelling dark matter contraction.

4.2 Shape and alignment of the dark matter haloes

Figure 2 shows our results for the shapes (see Section 2) and alignments of the dark matter and stellar mass distributions, while the values are listed in Table 3. In this Figure, the full distribution up to one pixel past the last-image radius R_L is probed. There are several key points to note. Firstly, the reconstructed dark matter distribution of all lenses considered here is rounder than the stellar distribution. Secondly, most lenses lie in the bottom-left corner: the dark matter distributions have typically low eccentricities, while the stellar distributions can be rather elliptical. Thirdly, 3 out of the 11 systems, 0712, 2016, and 1608, stand out due to their larger stellar ellipticity ($e_* > 0.25$). At the same time, they show a very strong alignment between their stellar and dark matter distributions. One of these lenses (1608) is a known merger of an early-type galaxy with a probable late-type galaxy (Surpi & Blandford 2003). We speculate that the other two systems, 0712 and 2016, may also be post-merger systems, explaining their similarities with 1608.

Fourthly, notice that all dark matter and stellar distributions with weak external shear ($\gamma < 0.1$) are aligned apart from 0047 which is almost spherical ($e_{dm} \sim 0.02$). By

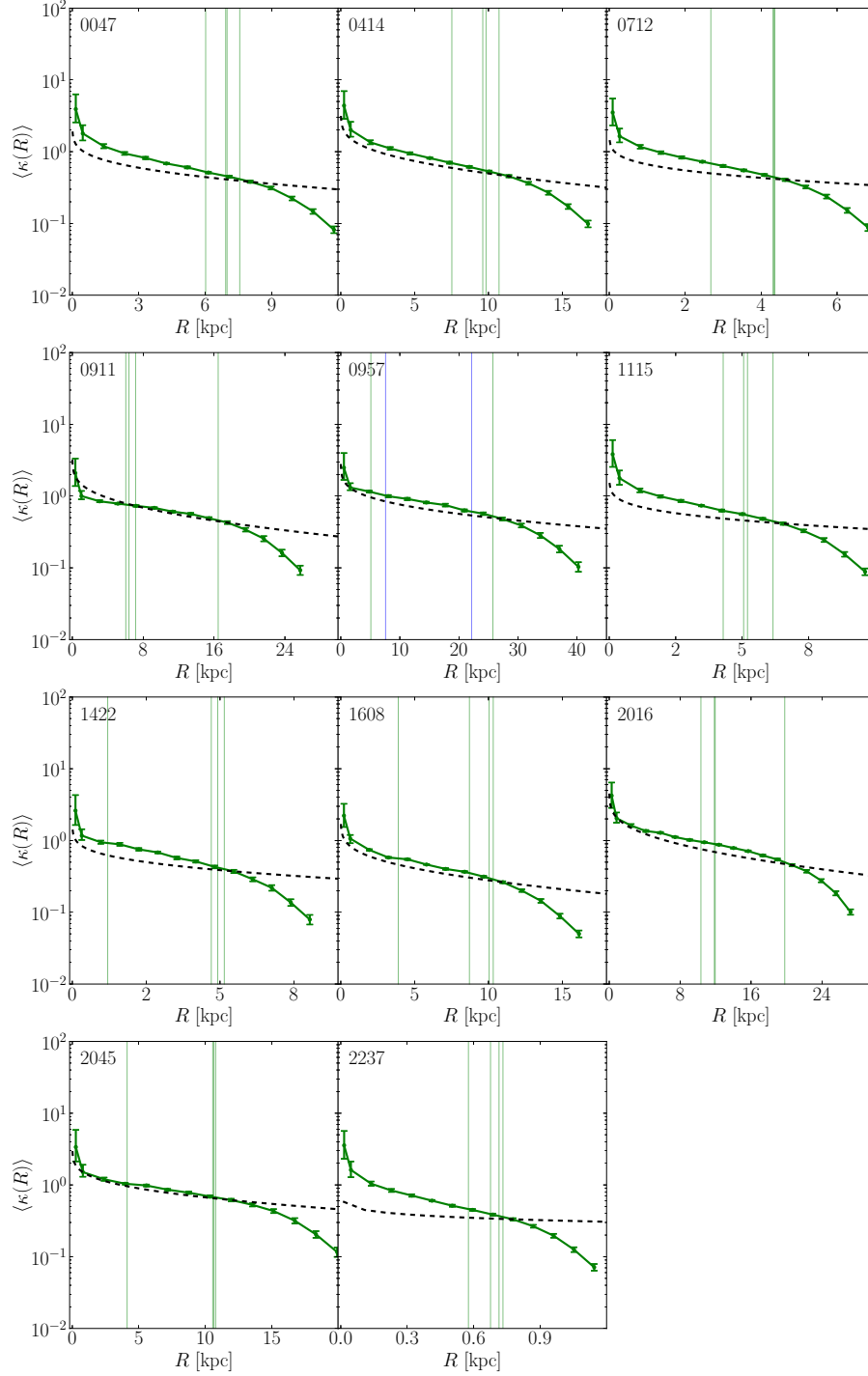


Figure 1. Azimuthally averaged convergence $\langle \kappa \rangle$ within radii R of the dark matter distribution. The vertical lines denote the radial distances of the source images from the centre of the lens galaxies. For 0957, there are two lensed components of the source galaxy denoted by the green and blue vertical lines. The green curves show the contributions by the dark matter in the strong lens galaxies as reconstructed by GLASS. The dashed black lines show the corresponding profiles for dark matter haloes that follow a NFW density profile. The parameters in the NFW distribution are calibrated such that the convergence $\langle \kappa \rangle$ matches the corresponding values of the reconstructed profiles at one pixel past the last image radius. Notice that those lenses where the outermost image is within ~ 10 kpc appear contracted as compared to an NFW profile; those where the outermost image is further out > 10 kpc are well fit by the NFW form. In all cases there is a sharp fall off in the GLASS dark matter profile at large radii due to the edge of the mass map.

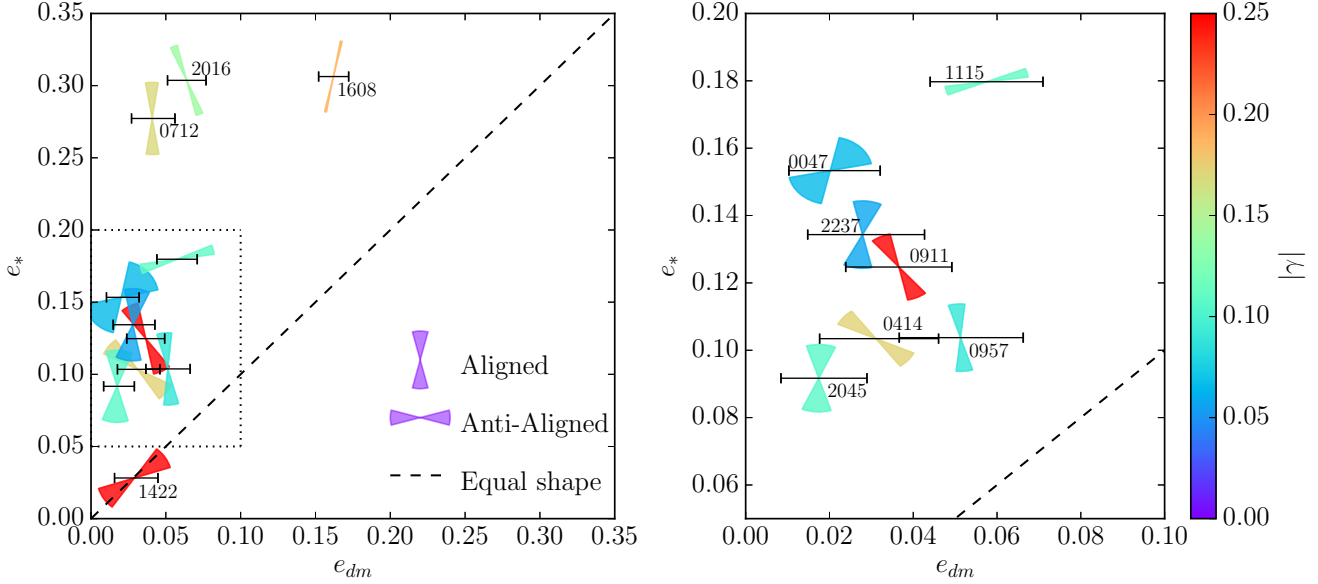


Figure 2. The reconstructed lens galaxies are displayed according to the reconstructed total shapes of the dark matter haloes e_{dm} and the stellar components e_* . The wedges show the alignment of the semi-major axes of the components. In case of alignment of the respective semi-major axes (Aligned), the wedges are vertical. In case of alignment of the semi-major with the semi-minor axis (Anti-Aligned), the wedges are horizontal. The opening angle of each wedge displays the statistical error in the shape estimate of the dark matter haloes. Both the uncertainties in the alignment and the ellipticity of the dark matter distribution denote the ranges 68% of all the mass models of the ensemble of solutions lie within. The lens galaxies are additionally color-coded in terms of the required external shear in reconstructing the mass distribution. The right panel shows the region in the left panel denoted by a dashed rectangle in greater detail.

Lens	Env	θ_{Env} [°]	r_{vir} [kpc]	M_{vir} [$10^{12} M_\odot$]	$ \gamma $	θ_g [°]	$\Delta\theta$ [°]	e_*	e_{dm}	e_{dm}/e_*	f_h
0047	G(9)	...	360.7	10.0	0.07 ± 0.01	120 ± 5	-41 ± 33	0.15	0.02 ± 0.01	0.13 ± 0.07	0.02 ± 0.15
0414	256.9	12.4	0.17 ± 0.02	166 ± 4	50 ± 15	0.10	0.03 ± 0.01	0.30 ± 0.14	-0.05 ± 0.15
0712	446.5	15.5	0.17 ± 0.01	141 ± 3	-0 ± 10	0.28	0.04 ± 0.01	0.15 ± 0.05	0.15 ± 0.05
0911	C	-160	369.1	22.7	0.28 ± 0.02	99 ± 2	25 ± 14	0.12	0.04 ± 0.01	0.29 ± 0.10	0.19 ± 0.12
0957	C	59	1053.1	175.2	0.09 ± 0.03	159 ± 8	4 ± 11	0.10	0.05 ± 0.01	0.49 ± 0.14	0.49 ± 0.14
1115	G(13)	-127	616.8	31.1	0.11 ± 0.01	154 ± 3	-75 ± 7	0.18	0.06 ± 0.01	0.32 ± 0.07	-0.28 ± 0.08
1422	G(17)	147	450.8	13.0	0.26 ± 0.01	39 ± 1	-51 ± 18	0.03	0.03 ± 0.01	1.03 ± 0.51	-0.23 ± 0.65
1608	G(8)	30 resp. -5^a	317.1	10.0	0.19 ± 0.01	145 ± 2	-12 ± 1	0.31	0.16 ± 0.01	0.53 ± 0.03	0.48 ± 0.03
2016	C(69)	130 – 140	299.3	22.1	0.14 ± 0.02	42 ± 4	20 ± 6	0.30	0.06 ± 0.01	0.21 ± 0.04	0.16 ± 0.04
2045	...	-60	562.8	104.0	0.11 ± 0.02	23 ± 7	-4 ± 20	0.09	0.02 ± 0.01	0.19 ± 0.11	0.19 ± 0.11
2237	1961.0	485.3	0.06 ± 0.01	159 ± 4	-7 ± 19	0.13	0.03 ± 0.01	0.21 ± 0.10	0.20 ± 0.11

Table 3. Fitted virial masses and radii and the ellipticity values of the stellar and dark matter components of the reconstructed strong lens galaxies and their corresponding ratios. ‘Env’ and θ_{Env} are the ‘Environment’ resp. θ_{Env} columns in Table 1; r_{vir} and M_{vir} are the virial radii and masses of the fitted NFW-profiles (see Section 4.1); $|\gamma|$ and θ_g denote magnitude and induced position angle (measured north through east) of the required external shear in the modeling of the strong lens galaxy; the ellipticities e_* and e_{dm} are defined by Eq. 2; $\Delta\theta = \theta_{dm} - \theta_*$ refers to the misalignment angle between the distributions of luminous and dark matter; e_{dm}/e_* denotes the ratio of the ellipticities of the dark matter relative to the stellar distribution; and $f_h = (e_{dm}/e_*) \cdot \cos 2\Delta\theta$ refers to the ratio of the ellipticity component of the dark matter distribution projected along the stellar distribution.

contrast, systems requiring a larger external shear ($\gamma > 0.1$; 0414, 0911, and 1422) can display a sizeable misalignment between the distributions. The latter two systems are located in dense environments (see Table 1) that are likely responsible for these misalignments. As described furthermore in Section 3, 1422 is a challenging system to model the stellar distribution. For 0414 there is another galaxy along the line-of-sight contributing to the lensing effect (Curran et al. 2011). Besides these three lens systems, one stands out from the rest, however. In 1115, the semi-major axis of the dark

matter seems to be aligned with the semi-minor axis of the stellar component; it is the most strongly misaligned (anti-aligned) lens system of all lenses studied here. This may be explained by its external shear ($\gamma \sim 0.1$) emanating from its group environment (see Table 1).

We performed tests for the robustness of our shape estimate and its dependence, in particular, on the shear prior. Only two systems, 1422 and 2045, reacted to a stricter prior on the external shear $|\gamma|$. 1422 requires a rather large external shear ($|\gamma| \sim 0.25$). Restricting the range of the allowed

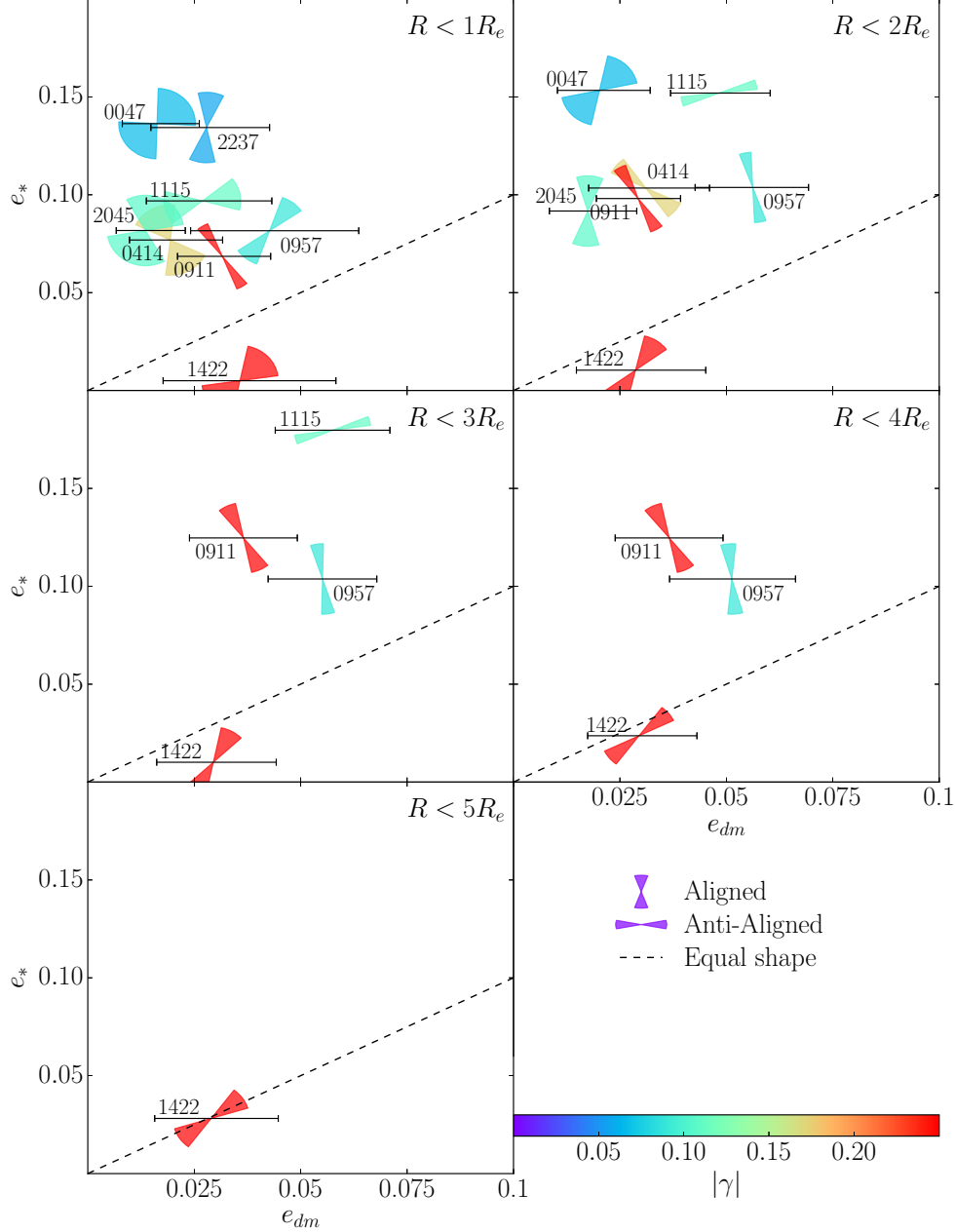


Figure 3. Similar to Figure 2. In each of the subfigures however we only consider pixels within 1, 2, 3, 4, or 5 R_e . If the outermost image of a system lies within one of the limiting radii, the system is kept in this panel, but dropped from subsequent panels.

external shear values did not affect the misalignment of the dark matter and stellar mass distributions. It did, however, increase the ellipticity of the reconstructed mass distribution: we observed a trade-off between the galaxy’s ellipticity and the magnitude of the external shear. *2045* on the other hand is particularly interesting. Allowing the shear to roam free as for the other lenses, favours an anomalously high shear ($\gamma > 0.5$), with strong misalignment. However, such models also produce spurious extra images along the arc. By limiting the maximum shear $|\gamma|$ to 0.1 the extra images are removed, though the lens then pushes on its maximum

shear prior. We find that the mass map is robust against a stronger shear prior, and does not affect the alignment of the distributions. It produces well-aligned dark matter and stellar distributions. We believe that a strong shear prior for this particular lens is justified by the appearance of spurious extra images if higher shear is allowed. We furthermore explore potential degeneracies between the estimation of the shape and the required external shear in the Appendix. As shown in Figure B1 the degeneracy can be largely broken for most lenses.

Figure 3 is similar to Figure 2 with the difference that

shapes are now averaged within different multiples of R_e . As the inertia tensor defined in Eq. 1 scales depends on the square of the radius, averaging up to different limiting radii could potentially yield different results. We observe that the shape estimator is robust to these changes. However, we note that there appears to be a slight trend that the ellipticity of the stellar distribution increases with increasing radius of enclosure. Whether this trend persists with data of a higher quality remains to be seen. There is potentially a similar trend for the dark matter distribution, although the samples are consistent with a constant ellipticity as well. 0957, and to some degree 1422, display, however, some signs of isophotal twist complicating the interpretation (see Appendix B).

5 DISCUSSION

5.1 Consistency with weak lensing estimates

In contrast to strong lensing, weak lensing as a probe cannot constrain e_{dm} directly. Rather, it is sensitive to the ratio $f_h = (e_{dm}/e_*) \cdot \cos 2\Delta\theta$, where $\Delta\theta$ is the misalignment angle between the two distributions, and $e_{dm} \cdot \cos 2\Delta\theta$ the component of the dark matter ellipticity projected along the light distribution. In case of perfect alignment and identical shape, f_h is equal to 1. In practice however, f_h deviates from 1, as among other effects misalignment of the two distributions and different ellipticities, which cannot be disentangled, change its value.

We compare our results presented in the previous Section with the most recent study of the shape and alignment of the dark matter halo relative to the distribution of luminous matter by Schrabback et al. (2015). They perform a measurement on CFHTLenS data and on data from the Millennium Simulation (Springel et al. 2005) with a ray-tracing through by Hilbert et al. (2009). The constraints on the early-type galaxies in the highest mass bin ($\log_{10} M_* > 11$) on CFHTLenS data are $f_h = -0.04^{+0.25}_{-0.25}$. On the simulated data, including cosmic shear and misalignments between the luminous and the dark matter, the constraints in the highest mass bin and for redshifts $0.4 < z < 0.6$ are $f_h = 0.359^{+0.011}_{-0.010}$. Due to a narrower intrinsic ellipticity distribution in the simulated data however, the latter constraints need to be corrected for with a factor $\sim 1/1.46$. Our results are broadly consistent with both constraints, with the exception of the galaxy merger system 1608 and the interesting and highly misaligned lens 1115. Our sample is, however, not large enough for a quantitative comparison.

5.2 The lack of isolated misaligned lenses

Our lens sample is small – just 11 strong lensing systems – yet it is interesting that we find no lens with strong misalignment and weak external shear ($\gamma < 0.1$). This is perhaps to be expected. Heiligman & Schwarzschild (1979) were the first to study the stability of orbits within triaxial figures, likely the physical situation in real elliptical galaxies. They found that while stable periodic orbits exist in all three symmetry planes of a triaxial figure, tube orbits about the intermediate axis are unstable. de Zeeuw (1985) studied the case of orbits misaligned to all three symmetry planes. They found that while misaligned stable orbits can

be found, they exist only for limited ranges of the Hamiltonian. Martinet & de Zeeuw (1988) extended these results to rotating triaxial figures showing that once rotation is included such orbits become even rarer. These early results are supported by more recent numerical work. Debattista et al. (2013) show that discs can never form perpendicular to the intermediate axis, in good agreement with earlier work by Heiligman & Schwarzschild (1979). However, they find that discs can survive off-axis if gas is present, since gas cooling can maintain the misalignment. Debattista et al. (2015b) go on to show that if this gas supply is switched off (as is the case for giant elliptical galaxies), then continuing collisionless minor mergers/interactions bring the disc rapidly into alignment with its host halo. Taken together, these results suggest that alignment between light and dark should be the norm for isolated gas-free triaxial galaxies, where here we define ‘isolated’ as having external shear $\gamma < 0.1$.

6 CONCLUSIONS

We have measured the projected radial density profile, shape and alignment of the stellar and dark matter distributions in 11 lens systems using a new non-parametric lens tool, GLASS. We focussed on lenses that have either time delay data or stellar mass maps that contribute significantly to the central potential. In a previous paper, we showed that data of this quality are required to determine the projected shape of dark matter haloes (Coles et al. 2014). We measured the shape and alignment using the eigenvalues λ_i and eigenvectors of the 2D moment of inertia tensor of the stellar and dark matter distributions, defining an estimator of the ellipticity $e = (\lambda_1 - \lambda_2)/(\lambda_1 + \lambda_2)$ ($\lambda_1 \geq \lambda_2$; see Section 2). We averaged e over the range $1-5R_e$, where R_e is the effective radius of the light profile.

Our key results are as follows:

- The projected dark matter density profile in these systems – under the assumption of a Chabrier stellar initial mass function – shows significant variation from galaxy to galaxy. Those with an outermost image beyond ~ 10 kpc are very well fit by a projected NFW profile; those with images within 10 kpc appear to be more concentrated than NFW, as expected if their dark haloes contract due to baryonic cooling. There could be several explanations for this behaviour. It could be a selection effect where lower mass galaxies (expected to have closer-in images) can only strong lens if they are substantially contracted; it could owe to more massive haloes (with further out images) being less affected by baryons; or it could be that we simply cannot yet resolve the contraction at the centre of the more massive systems. We will explore such questions in a forthcoming paper.
- The dark matter haloes of all lenses in our sample are rounder than the light distribution over the range $1R_e < R < 5R_e$. As we average over larger radii, there is a slight trend for the lenses, except the ones with a large stellar ellipticity, to become increasingly elliptical in their stellar distributions. A similar behaviour can be seen also for the dark matter, though the effect is not statistically significant. The dark matter haloes are never more elliptical than $e_{dm} = 0.2$, while their stars can extend to $e_* > 0.2$.
- Three systems have a high stellar ellipticity ($e_* > 0.25$) and correspondingly high alignment between light and dark.

One of these – 1608 – is a known merging pair. We speculate that the other two lens systems (0712 and 2016) may also be recent post-merger systems.

- Galaxies with high dark matter ellipticity and weak external shear ($\gamma < 0.1$) show strong alignment; those with strong shear ($\gamma \gtrsim 0.1$) can be highly misaligned. This is reassuring since isolated misaligned galaxies are expected to be unstable (e.g. Heiligman & Schwarzschild 1979; Martinet & de Zeeuw 1988; Adams et al. 2007; Debattista et al. 2015a). We find that lenses with external shear $\gamma < 0.1$ appear to be sufficiently isolated that their luminous and dark matter distributions are well-aligned.

Our results on the shape and radial profile of the dark matter haloes in these lenses provide new constraints on galaxy formation models. Dark matter haloes can contract due to baryonic cooling (e.g. Blumenthal et al. 1986) or expand due to energetic feedback from supernovae (e.g. Read & Gilmore 2005; Pontzen & Governato 2012, 2014; Oñorbe et al. 2015; Read et al. 2015) or AGN (e.g. Martizzi et al. 2012). The latest simulations for elliptical galaxies forming in Λ CDM without AGN feedback suggest that contraction should win (Dutton et al. 2015). Under the assumption of a Chabrier IMF, our results imply that AGN do not play a major role in the lower mass lenses in our sample (since these are contracted with respect to NFW), but perhaps do act to counterbalance contraction in the more massive lenses. For the shape and alignment of lens galaxies, the models must reproduce rather round dark matter halos within $5R_e$, and misaligned stellar and dark matter distributions in lenses with strong external shear ($\gamma > 0.1$). It will be interesting to see whether the latest galaxy formation models reproduce this naturally, or it poses a challenge. Finally, we note that strong misalignments may be problematic for those alternative gravity theories in which light is the only source of gravity, since in such models light and dark must necessarily be highly correlated.

7 ACKNOWLEDGEMENTS

We thank Simon Birrer and William Hartley for useful discussion. We also thank the anonymous referee for constructive comments. JIR would like to acknowledge support from SNF grant PP00P2.128540/1. DL’s research is part of the project GLENCO, funded under the European Seventh Framework Programme, Ideas, Grant Agreement n. 259349.

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APPENDIX A: RECONSTRUCTED LENSES

In this Appendix, we show the results of our lens modelling for each individual lens galaxy (Figures A1-A3). The panels show, from left to right: the arrival time surface; the surface mass density of the dark matter; and the surface mass density of the stars. The solid lines mark the eigenvalues and eigenvectors of the 2D moment of inertia tensor in each case; the dashed lines denote the ranges 68% of all the eigenvectors of the dark matter distribution corresponding to an individual model lie within. Figures 2 & 3 are constructed from a combination of the eigenvalues (Eq. 2), and the angle between the dark matter and stellar major axes (Eq. 3). Note that the angular scale is always the same for the dark matter and stellar maps, but varies between the different lenses as marked on the Figure axes.

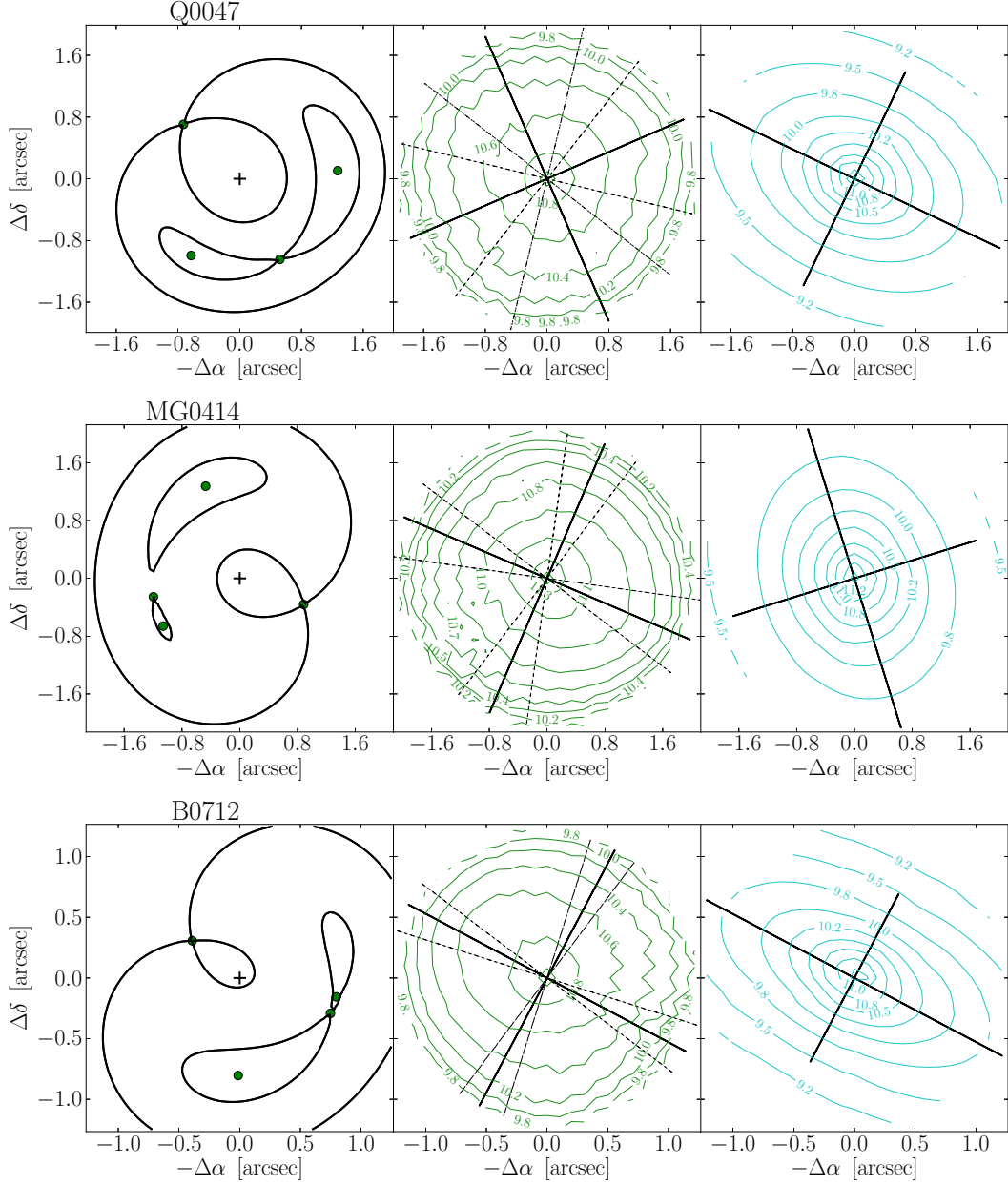


Figure A1. The results of our reconstruction of each individual lens galaxy. The panels show, from left to right: the arrival time surface (the images are marked by the green circles); the surface mass density of the dark matter; and the surface mass density of the stars. The solid lines mark the eigenvalues and eigenvectors of the 2D moment of inertia tensor in each case; the dashed lines the 68% confidence interval of these for the dark matter map.

We note in Figure A2, specifically for *0957* and to some degree also *1422*, twisting isophotes (e.g. Binney 1978).

APPENDIX B: DEGENERACY BETWEEN θ_{DM} AND θ_G

Figure B1 shows the degeneracy between the position angles of the dark matter halo θ_{dm} and the external shear θ_g for each strong lens galaxy. We note that the constraints on the direction of the external shear are stronger than on

the dark matter distribution. There is a weak degeneracy between the two quantities, it can however for most lenses be largely broken. This is expected, since internal ellipticity and external shear tend to have similar effects, but are not an exact degeneracy (cf. Dominik 1999).

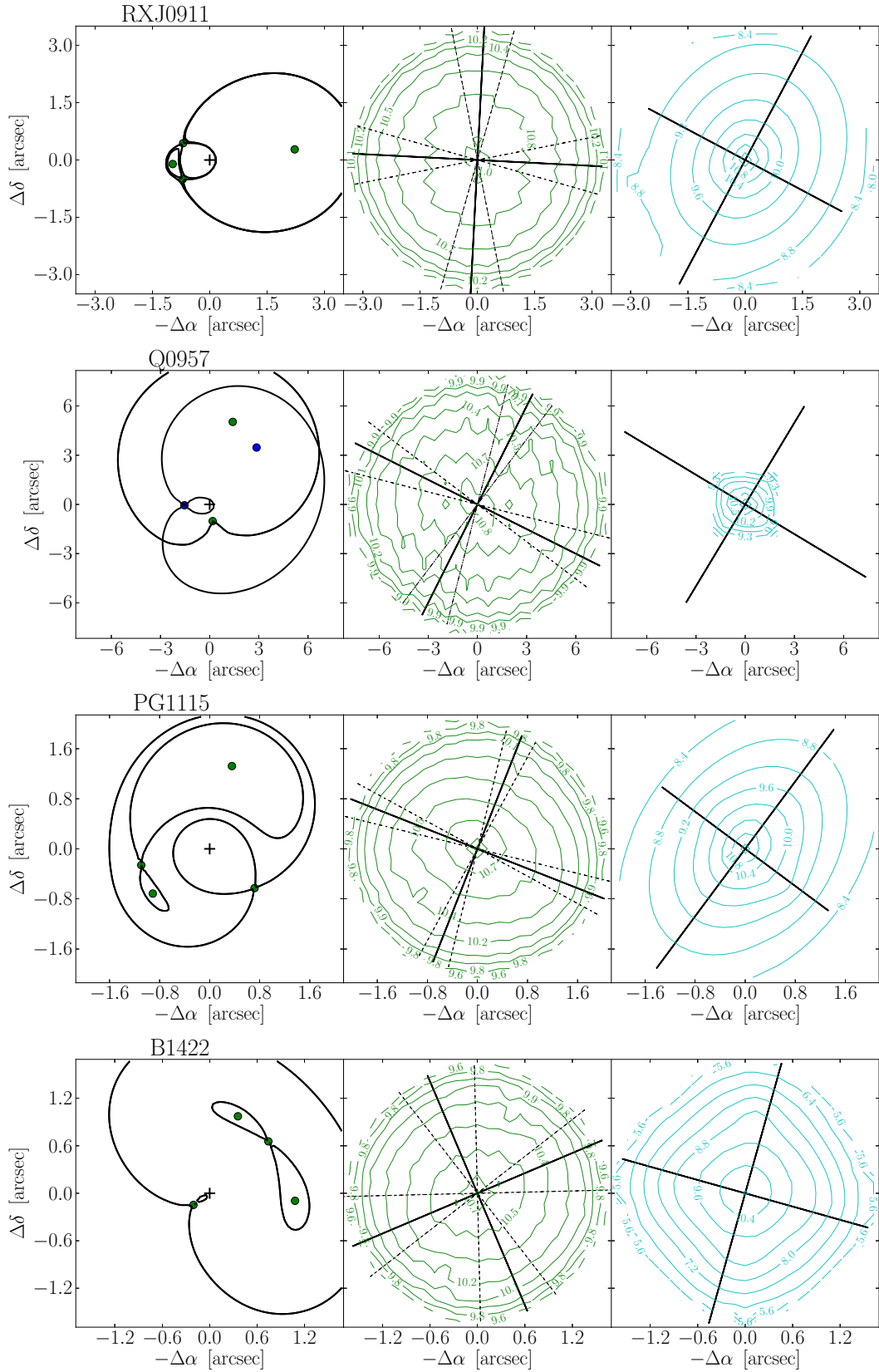


Figure A2. The results of our reconstruction of each individual lens galaxy. Lines and symbols are as in Figure A1.

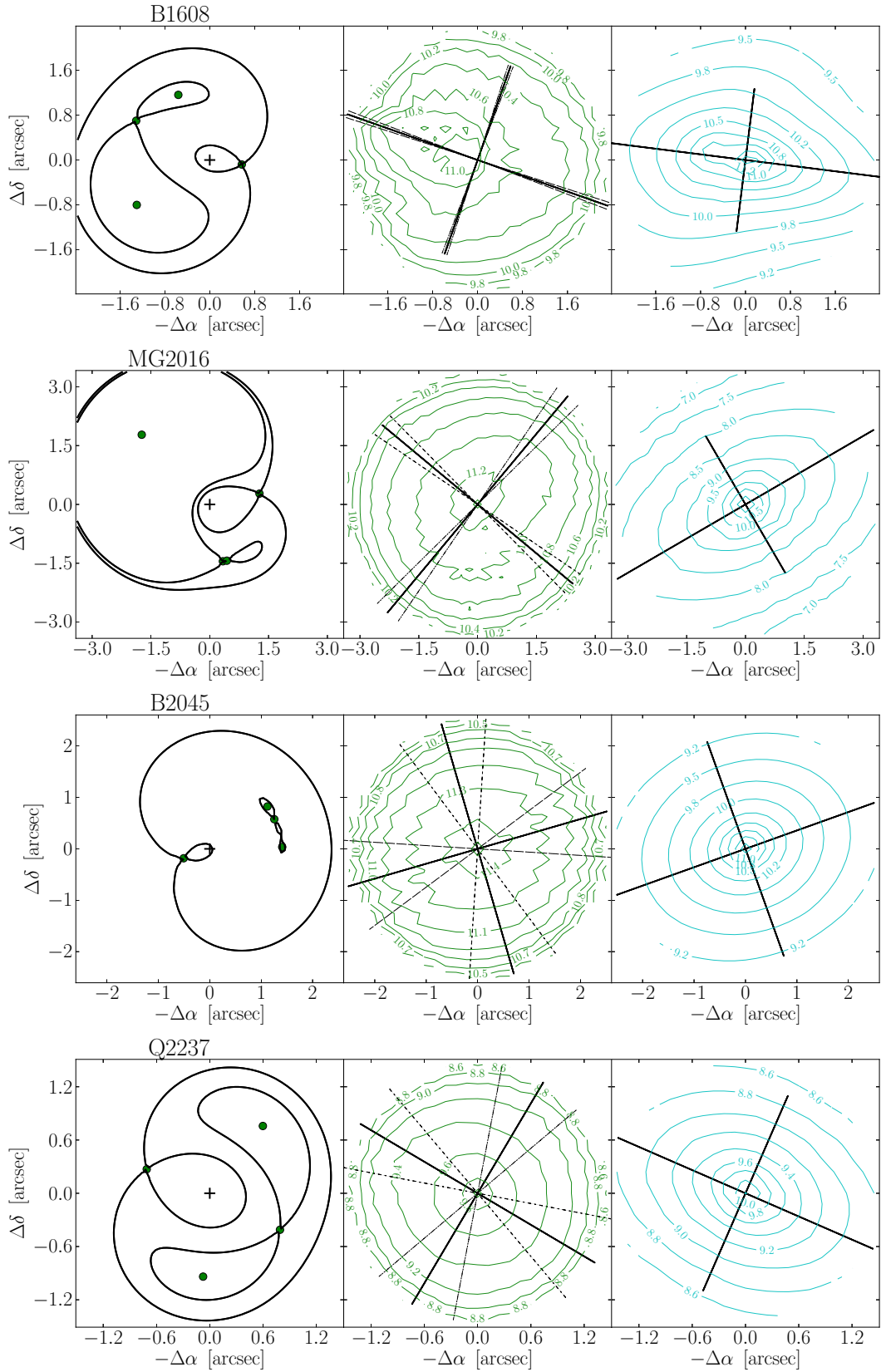


Figure A3. The results of our reconstruction of each individual lens galaxy. Lines and symbols are as in Figure A1.

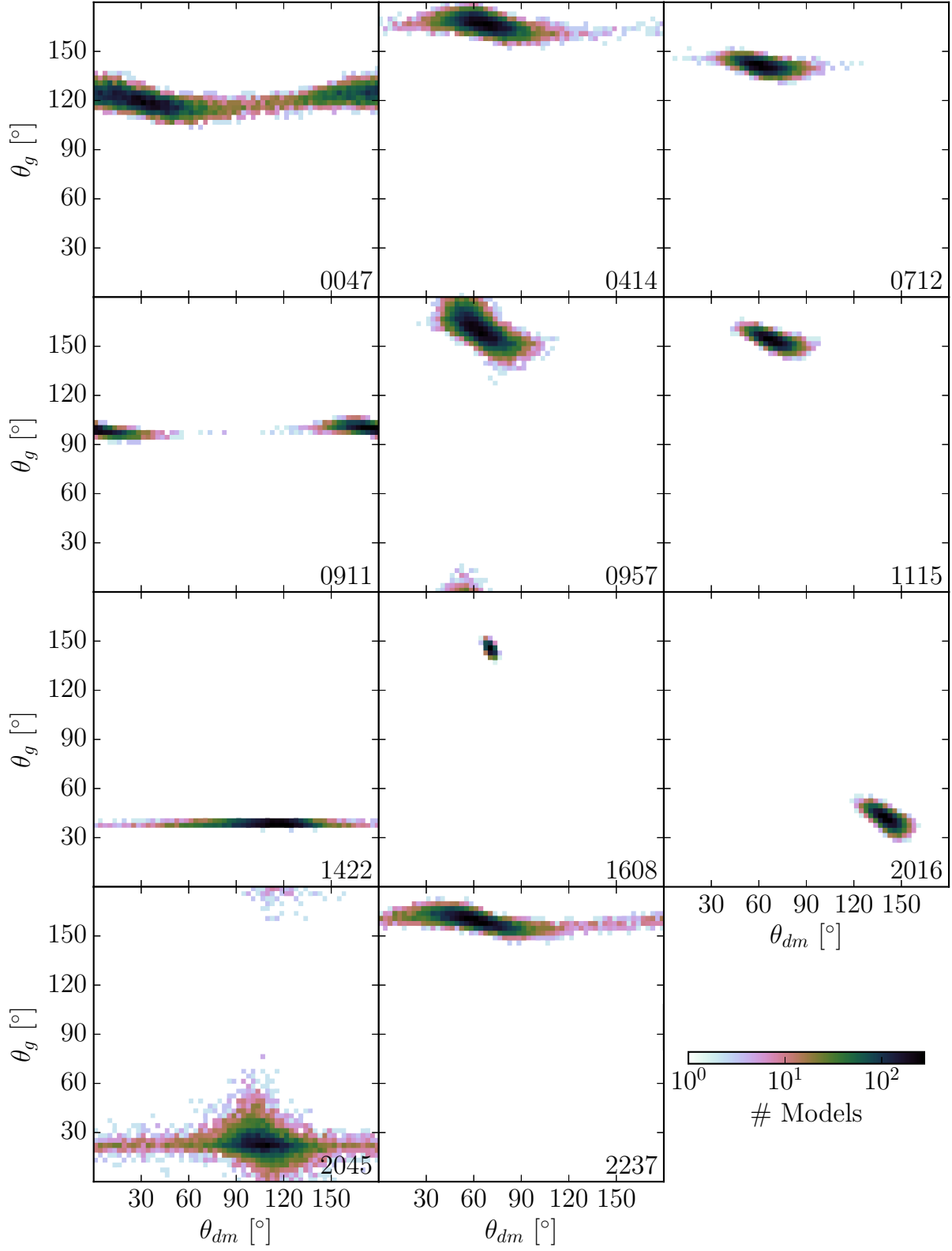


Figure B1. The two-dimensional distribution of the position angles of the dark matter distribution θ_{dm} and the induced position angle of the required external shear θ_g of each reconstructed mass distribution. The panels display this distribution, which traces the degeneracy of the two parameters, for each reconstructed strong lens galaxy in the ensemble of solutions. Here, the angles are rotated to be measured north through east.

Lens	A ["]	B ["]	C ["]	D ["]
0047	1.270	-0.630	0.520	-0.730
	0.105	-0.995	-1.045	0.705
0414	-0.472	-1.061	-1.1947	0.885
	1.277	-0.661	-0.255	-0.361
0712	-0.013	0.795	0.747	-0.391
	-0.804	-0.156	-0.292	0.307
0911	2.226	-0.968	-0.709	-0.696
	0.278	-0.105	-0.507	0.439
0957 ^a	1.408	0.182	2.860	-1.540
	5.034	-1.018	3.470	-0.050
1115	0.355	-0.909	-1.093	0.717
	1.322	-0.714	-0.260	-0.627
1422	1.079	0.357	0.742	-0.205
	-0.095	0.973	0.656	-0.147
1608	-1.300	-0.560	-1.310	0.570
	-0.800	1.160	0.700	-0.080
2016	-1.735	0.335	0.437	1.268
	1.778	-1.450	-1.435	0.276
2045	1.121	1.409	1.255	-0.507
	0.824	0.035	0.576	-0.183
2237	0.598	-0.075	0.791	-0.710
	0.758	-0.939	-0.411	0.271

Table A1. The image positions of all the systems are listed in order of arrival time (image A has the shortest arrival time, image D the longest). The positions are relative to the centre of the lens galaxy. The first row contains the RA coordinate, the second the DEC coordinate of the image.

^a The source lensed by 0957 has a second component. Images A and B are a double associated with the main galaxy component, images C and D with the lensed substructure.